Hyperbolcity-Preserving Well-Balanced Stochastic Galerkin Method for Shallow Water Equations

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Abstract

We study the stochastic Galerkin (SG) method for stochastic parameterized shallow water equations. Our work comprises the following aspects:

- A hyperbolicity-preserving stochastic Galerkin formulation for the shallow water equations using only the conserved variables.
- A sufficient condition to preserve the hyperbolicity, which is a stochastic variant of the deterministic positivity condition.
- A computationally tractable condition to guarantee the hyperbolicity.
- A central-upwind scheme that preserves both the hyperbolicity and the well-balanced property at discrete time levels.

Motivations

- Uncertainties can enter the shallow water system, for example, via the noisy measurement of the bottom.
- A SG formulation of shallow water equations is not necessarily hyperbolic.
- A non-well-balanced scheme may lead to spurious oscillations on relatively coarse grid.

Stochastic Parameterized Shallow Water System

$$\frac{\partial h}{\partial t} + \frac{\partial q^x}{\partial x} + \frac{\partial q^y}{\partial y} = 0,$$

$$\frac{\partial q^x}{\partial t} + \frac{\partial}{\partial x} \left(\frac{(q^x)^2}{h} + \frac{gh^2}{2} \right) + \frac{\partial}{\partial y} \left(\frac{q^x q^y}{h} \right) = -gh \frac{\partial B}{\partial x},$$

$$\frac{\partial q^y}{\partial t} + \frac{\partial}{\partial x} \left(\frac{q^x q^y}{h} \right) + \frac{\partial}{\partial y} \left(\frac{(q^y)^2}{h} + \frac{gh^2}{2} \right) = -gh \frac{\partial B}{\partial y}.$$
h is the water height, a_x and a_y are the x - and y -

h is the water height, q_x and q_y are the x- and ydischarges, and B is the time-independent surface. All the variables are ξ -dependent random fields, e.g., $h = h(x, y, t, \xi).$

Polynomial Chaos Expansion (PCE)

• An unknown random field $z(x, y, t, \xi)$ is represented in the L^2_{ρ} orthonormal basis $\{\phi_k\}_{k\in\mathbb{R}}$, where $\rho := \rho(\xi)$ is the density of the random parameter.

$$z(x,y,t,\xi) = \sum_{k=1}^{\infty} \hat{z}_k(x,y,t)\phi_k(\xi),$$

• K-term truncated PCE:

$$\Pi_{\Lambda}[z]\coloneqq\sum_{k=1}^K\widehat{z}_k(x,y,t)\phi_k(\xi),$$

where Λ is the index set for the (possibly multivariate) polynomials, the cardinality of Λ is $K, \text{ and } \phi_1(\xi) = 1.$

• A K-term PCE approximation to product of two random fields a and b:

$$\Pi_{\Lambda}[a,b] \coloneqq \Pi_{\Lambda}[\Pi_{\Lambda}[a] \Pi_{\Lambda}[b]].$$

• A K-term PCE approximation to the ratio of two random fields a and b:

 $\Pi_{\Lambda}^{\dagger}[b/a]$: the solution to $\Pi_{\Lambda}[a,b/a]=\Pi_{\Lambda}[b]$.

Notations

- $\bullet \, \widehat{z} = (\widehat{z}_1, \dots, \widehat{z}_K)^{\top}.$
- $\mathcal{P}(\hat{z}) := \sum_{k=1}^K \hat{z}_k \mathcal{M}_k, \quad (\mathcal{M}_k)_{\ell m} = \langle \phi_k, \phi_\ell \phi_m \rangle_{\rho}.$
- It can be shown that

 $\widehat{\Pi_{\Lambda}[a,b]} = \mathcal{P}(\widehat{a})\widehat{b}, \qquad \widehat{\Pi_{\Lambda}^{\dagger}[b/a]} = \mathcal{P}^{-1}(\widehat{a})\widehat{b}.$

Stochastic Galerkin (SG) Method

• Ansatz:

Insatz.
$$h \simeq h_{\Lambda} \coloneqq \sum_{k=1}^{K} \widehat{h}_{k}(x,y,t) \phi_{k}(\xi),$$

$$q^{x} \simeq q_{\Lambda}^{x} \coloneqq \sum_{k=1}^{K} (\widehat{q^{x}})_{k}(x,y,t) \phi_{k}(\xi),$$

$$q^{y} \simeq q_{\Lambda}^{y} \coloneqq \sum_{k=1}^{K} (\widehat{q^{y}})_{k}(x,y,t) \phi_{k}(\xi),$$

• Stochastic Galerkin method applies standard Galerkin procedure in the stochastic ξ space, which leads to a new system of partial differential equations with respect to the PCE coefficients.

SG Projection of Nonlinear Terms

 $\frac{(q^x)^2}{h} = \frac{q^x}{h} q^x \longrightarrow \Pi_{\Lambda} \left[\frac{(q_{\Lambda}^x)^2}{h_{\Lambda}} \right] = \Pi_{\Lambda} \left[q_{\Lambda}^x \Pi_{\Lambda}^{\dagger} \left[\frac{q_{\Lambda}^x}{h_{\Lambda}} \right] \right], \qquad \eta(x, y, 0, \xi) = \begin{cases} 1.01, & \text{if } 0.05 < x < 0.15, \\ 1, & \text{otherwise,} \end{cases}$ $\frac{(q^y)^2}{h} = \frac{q^y}{h} \ q^y \longrightarrow \Pi_{\Lambda} \left[\frac{(q^y_{\Lambda})^2}{h_{\Lambda}} \right] = \Pi_{\Lambda} \left[q^y_{\Lambda} \ \Pi^{\dagger}_{\Lambda} \left[\frac{q^y_{\Lambda}}{h_{\Lambda}} \right] \right].$

• For $q^x q^y/h$ in $(q^x q^y/h)_x$,

$$\frac{q^x q^y}{h} = (q^x) \frac{q^y}{h} \longrightarrow \Pi_{\Lambda} \left[\frac{q_{\Lambda}^x q_{\Lambda}^y}{h_{\Lambda}} \right] = \Pi_{\Lambda} \left[q_{\Lambda}^x \Pi_{\Lambda}^{\dagger} \left[\frac{q_{\Lambda}^y}{h_{\Lambda}} \right] \right].$$

Sor $q^x q^y/h$ in $(q^x q^y/h)_y$,

$$\frac{q^x q^y}{h} = (q^y) \frac{q^x}{h} \longrightarrow \Pi_{\Lambda} \left[\frac{q_{\Lambda}^x q_{\Lambda}^y}{h_{\Lambda}} \right] = \Pi_{\Lambda} \left[q_{\Lambda}^y \Pi_{\Lambda}^{\dagger} \left[\frac{q_{\Lambda}^x}{h_{\Lambda}} \right] \right].$$

Numerical simulations

Deterministic initial water surface:

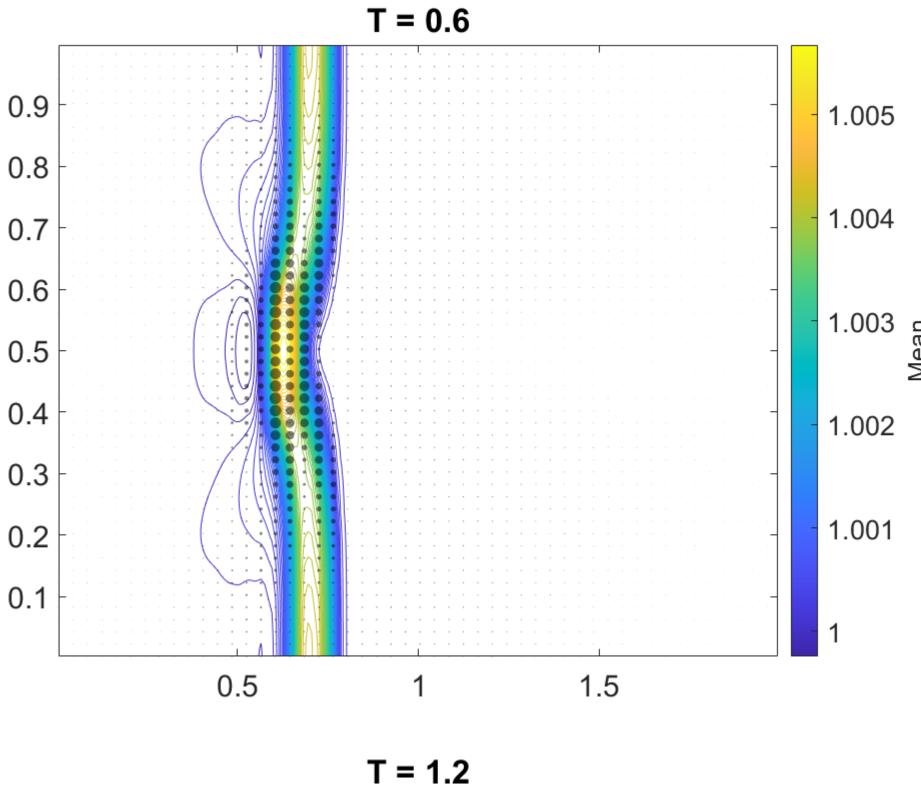
$$\eta(x, y, 0, \xi) = \begin{cases} 1.01, & \text{if } 0.05 < x < 0.15 \\ 1, & \text{otherwise,} \end{cases}$$

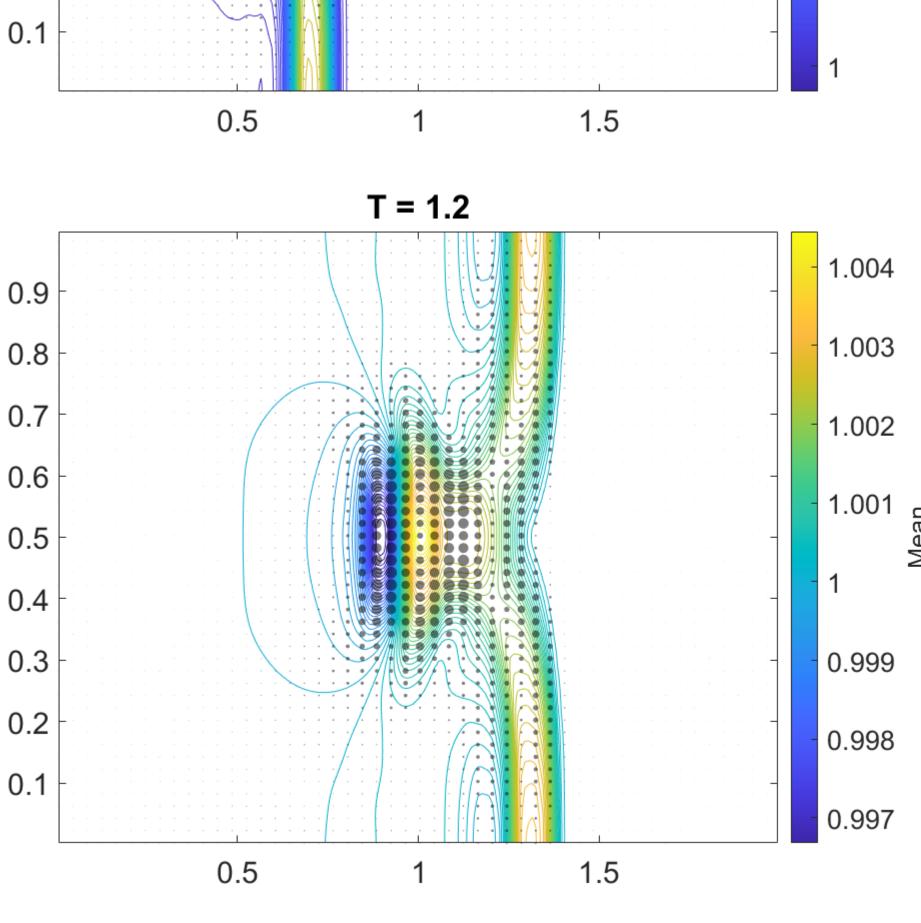
Deterministic initial velocity field (0-discharge): $u(x, y, 0, \xi) = v(x, y, 0, \xi) = 0.$

Stochastic bottom topography:

$$B(x, y, \xi) = 0.8e^{-5(x-0.9+0.1\xi^{(1)})^2 - 50(y-0.5+0.1\xi^{(2)})^2}$$
. Randomness:

 $\xi^{(1)} \sim \text{Beta}(4,2),$ $\xi^{(2)} \sim \mathcal{U}(-1,1).$





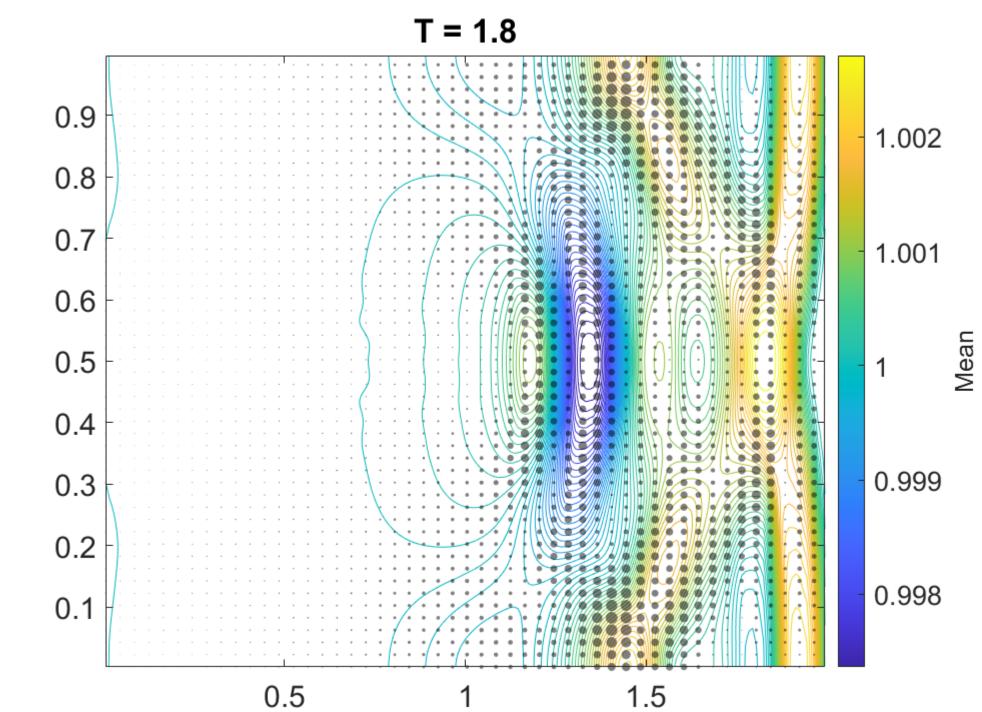


Figure 1:Numerical results at T = 0.6 (top), T = 1.2 (middle), and T = 1.8 (bottom), respectively. The largest disks are corresponding to the standard deviation values 2.20e-3, 2.00e-3, and 1.20e-3, respectively. The index set $\Lambda = \{(\nu^{(1)}, \nu^{(2)}) \in$ $\mathbb{N}^2 \mid 0 \leq \nu^{(1)}, \nu^{(2)} \leq 3$. The polynomial basis is chosen to be the tensor-product set. Animation can be found in https: //ibit.ly/Q6H4

$$\frac{\partial}{\partial t}(\widehat{U}) + \frac{\partial}{\partial x}(\widehat{F}(\widehat{U})) + \frac{\partial}{\partial y}(\widehat{G}(\widehat{U})) = \widehat{S}(\widehat{U}, \widehat{B}). \tag{1}$$
and

Here,
$$\widehat{U} := (\widehat{h}^{\top}, \widehat{q^x}^{\top}, \widehat{q^y}^{\top})^{\top}$$
, and
$$\widehat{F}(\widehat{U}) = \begin{pmatrix} \widehat{q^x} \\ \mathcal{P}(\widehat{q^x})\mathcal{P}^{-1}(\widehat{h})\widehat{q^x} + \frac{1}{2}g\mathcal{P}(\widehat{h})\widehat{h} \\ \mathcal{P}(\widehat{q^x})\mathcal{P}^{-1}(\widehat{h})\widehat{q^y} \end{pmatrix}, \quad \widehat{G}(\widehat{U}) = \begin{pmatrix} \widehat{q^y} \\ \mathcal{P}(\widehat{q^y})\mathcal{P}^{-1}(\widehat{h})\widehat{q^x} \\ \mathcal{P}(\widehat{q^y})\mathcal{P}^{-1}(\widehat{h})\widehat{q^y} + \frac{1}{2}g\mathcal{P}(\widehat{h})\widehat{h} \end{pmatrix}, \quad \widehat{S}(\widehat{U}, \widehat{B}) = \begin{pmatrix} 0 \\ -g\mathcal{P}(\widehat{h})\widehat{B_x} \\ -g\mathcal{P}(\widehat{h})\widehat{B_y} \end{pmatrix}. \quad (2)$$

The Main Results

Hyperbolicity-Preserving SG Formulation

Theorem (Hyperbolicity-preserving condition)

The system (1) is hyperbolic if the matrix $\mathcal{P}(\hat{h}) > 0$.

The condition $\mathcal{P}(\hat{h}) > 0$ reduces to h > 0 when the ξ -dependence is dropped from the system.

Theorem (A computationally tractable condition)

Given Λ , let nodes ξ_m and weights τ_m satisfying $\{(\xi_m, \tau_m)\}_{m=1}^M$ represent any M-point positive quadrature rule that is exact on

$$P_{\Lambda}^3 \coloneqq \operatorname{span} \left\{ \prod_{n=1}^3 \phi_n \mid n \in [K] \right\}.$$

$$h_{\Lambda}(x, y, t, \xi_m) > 0 \quad \forall \ m = 1, \dots, M,$$

then the matrix $\mathcal{P}(\hat{h}) > 0$.

In other words, we only need to ensure the positivity of the stochastic water heights at some quadrature points to preserve the hyperbolicity of (1).

Second-Order Central-Upwind Scheme

Assuming uniform rectangular partition over a rectangular region,

$$\frac{d}{dt}\boldsymbol{U}_{i,j} = -\frac{\boldsymbol{\mathcal{F}}_{i+\frac{1}{2},j} - \boldsymbol{\mathcal{F}}_{i-\frac{1}{2},j}}{\Delta x} - \frac{\boldsymbol{\mathcal{G}}_{i,j+\frac{1}{2}} - \boldsymbol{\mathcal{G}}_{i,j-\frac{1}{2}}}{\Delta y} + \overline{\boldsymbol{S}}_{i,j}$$

where $\boldsymbol{U}_{i,j}$ represent the cell averages of the vector \widehat{U} in rectangular cell $\mathcal{C}_{i,j}$.

• Source term:

$$\overline{m{S}}_{i,j} pprox rac{1}{|\mathcal{C}_{i,j}|} \int_{\mathcal{C}_{i,j}} \widehat{S}(\widehat{U},\widehat{B}) dx dy.$$

• Numerical fluxes:

Numerical fluxes:
$$\mathcal{F}_{i+\frac{1}{2},j} \coloneqq \frac{a_{i+\frac{1}{2},j}^{+} \widehat{F}(\boldsymbol{U}_{i,j}^{E}) - a_{i+\frac{1}{2},j}^{-} \widehat{F}(\boldsymbol{U}_{i+1,j}^{W})}{a_{i+\frac{1}{2},j}^{+} - a_{i+\frac{1}{2},j}^{-}} + \frac{a_{i+\frac{1}{2},j}^{+} a_{i+\frac{1}{2},j}^{-} a_{i+\frac{1}{2},j}^{-}}{a_{i+\frac{1}{2},j}^{+} - a_{i+\frac{1}{2},j}^{-}} \left[\boldsymbol{U}_{i+1,j}^{W} - \boldsymbol{U}_{i,j}^{E} \right]$$

 $\mathcal{G}_{i+rac{1}{2},j} \coloneqq rac{b_{i,j+rac{1}{2}}^{+}\widehat{G}(oldsymbol{U}_{i,j}^{N}) - b_{i,j+rac{1}{2}}^{-}\widehat{G}(oldsymbol{U}_{i,j+rac{1}{2}}^{S})}{b_{i,j+rac{1}{2}}^{+} - b_{i,j+rac{1}{2}}^{-}} + rac{b_{i,j+rac{1}{2}}^{+}b_{i,j+rac{1}{2}}^{-}}{b_{i,j+rac{1}{2}}^{+} - b_{i,j+rac{1}{2}}^{-}} \left[oldsymbol{U}_{i,j+1}^{S} - oldsymbol{U}_{i,j}^{N}
ight].$

• Propagation speeds:

$$a_{i+\frac{1}{2},j}^{-} = \min \left\{ \lambda_{1} \left(\frac{\partial \widehat{F}}{\partial \widehat{U}} (\boldsymbol{U}_{i+1,j}^{W}) \right), \lambda_{1} \left(\frac{\partial \widehat{F}}{\partial \widehat{U}} (\boldsymbol{U}_{i,j}^{E}) \right), 0 \right\},$$

$$a_{i+\frac{1}{2},j}^{+} = \max \left\{ \lambda_{3K} \left(\frac{\partial \widehat{F}}{\partial \widehat{U}} (\boldsymbol{U}_{i+1,j}^{W}) \right), \lambda_{3K} \left(\frac{\partial \widehat{F}}{\partial \widehat{U}} (\boldsymbol{U}_{i,j}^{E}) \right), 0 \right\},$$

$$b_{i,j+\frac{1}{2}}^{-} = \min \left\{ \lambda_{1} \left(\frac{\partial \widehat{G}}{\partial \widehat{U}} (\boldsymbol{U}_{i,j+1}^{S}) \right), \lambda_{1} \left(\frac{\partial \widehat{G}}{\partial \widehat{U}} (\boldsymbol{U}_{i,j}^{N}) \right), 0 \right\},$$

$$b_{i,j+\frac{1}{2}}^{+} = \max \left\{ \lambda_{3K} \left(\frac{\partial \widehat{G}}{\partial \widehat{U}} (\boldsymbol{U}_{i,j+1}^{S}) \right), \lambda_{3K} \left(\frac{\partial \widehat{G}}{\partial \widehat{U}} (\boldsymbol{U}_{i,j}^{N}) \right), 0 \right\},$$

 $U_{i,i}^{E,W,N,S}$ are the pointwise values of the secondorder accurate, non-oscillatory piecewise linear reconstructions of $\boldsymbol{U}_{i,j}$ at the midpoints of the boundaries, i.e.,

$$egin{aligned} oldsymbol{U}_{i,j}^E &= oldsymbol{U}_{i,j} + rac{\Delta x}{2} (oldsymbol{U}_x)_{i,j}, & oldsymbol{U}_{i,j}^W &= oldsymbol{U}_{i,j} - rac{\Delta x}{2} (oldsymbol{U}_x)_{i,j}, \ oldsymbol{U}_{i,j}^N &= oldsymbol{U}_{i,j} + rac{\Delta y}{2} (oldsymbol{U}_y)_{i,j}, & oldsymbol{U}_{i,j}^S &= oldsymbol{U}_{i,j} - rac{\Delta y}{2} (oldsymbol{U}_y)_{i,j}, \end{aligned}$$

Hyperbolicity-Preserving Well-Balanced Central-Upwind Scheme

• Stochastic "lake-at-rest" state:

$$\begin{cases} q_{\Lambda}^{x} = q_{\Lambda}^{y} \equiv 0, \\ h_{\Lambda} + \Pi_{\Lambda}[B] \equiv C(\xi), \end{cases} \Rightarrow \begin{cases} \widehat{q^{x}} = \widehat{q^{y}} \equiv \mathbf{0}, \\ \widehat{h} + \widehat{B} \equiv \widehat{C}. \end{cases}$$

- The PCE vector \widehat{B} for the bottom function is replaced by its piecewise bilinear interpolant.
- The pointwise values of the reconstructions of the PCE of water surface $\hat{\eta}$ are reconstructed. The reconstructed water height are computed by $\hat{h} \coloneqq \hat{\eta} - \hat{B}.$
- The first moments \hat{h}_1 are "corrected" following a similar procedure to the central-upwind scheme for the deterministic shallow water equations.
- The PCE vectors \hat{h} are filtered to satisfies the condition (3).

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