Hyperbolicity-Preserving Well-Balanced Stochastic Galerkin Method for Shallow Water Equations

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Abstract	Stochastic Galerkin (SG) Method	SG Projection of Nonlinear Terms	Numerical simulations
We study the stochastic Galerkin (SG) method for stochastic parameterized shallow water equa- tions. Our work comprises the following aspects:	• Ansatz: $h \simeq h_{\Lambda} \coloneqq \sum_{\substack{k=1 \ K}}^{K} \hat{h}_k(x, y, t) \phi_k(\xi),$	$\frac{(q^{x})^{2}}{h} = \frac{q^{x}}{h} q^{x} \longrightarrow \Pi_{\Lambda} \left[\frac{(q^{x}_{\Lambda})^{2}}{h_{\Lambda}} \right] = \Pi_{\Lambda} \left[q^{x}_{\Lambda} \Pi^{\dagger}_{\Lambda} \left[\frac{q^{x}_{\Lambda}}{h_{\Lambda}} \right] \right],$	Deterministic initial water surface: $\eta(x, y, 0, \xi) = \begin{cases} 1.01, \text{ if } 0.05 < x < 0.15, \\ 1, & \text{otherwise,} \end{cases}$
• A hyperbolicity-preserving stochastic Galerkin formulation for the shallow water equations	$egin{aligned} q^x &\simeq q^x_\Lambda \coloneqq \sum_{k=1}^{\kappa} (\widehat{q^x})_k(x,y,t) \phi_k(\xi), \ q^y &\simeq q^y_\Lambda \coloneqq \sum_{k=1}^{K} (\widehat{q^y})_k(x,y,t) \phi_k(\xi), \end{aligned}$	$\frac{(q^{s})^{2}}{h} = \frac{q^{s}}{h} q^{y} \longrightarrow \Pi_{\Lambda} \left[\frac{(q^{s}_{\Lambda})^{2}}{h_{\Lambda}} \right] = \Pi_{\Lambda} \left[q^{y}_{\Lambda} \Pi^{\dagger}_{\Lambda} \left[\frac{q^{s}_{\Lambda}}{h_{\Lambda}} \right] \right].$ $\text{2 For } q^{x}q^{y}/h \text{ in } (q^{x}q^{y}/h)_{x},$ $x = u = u = \left[\int_{X} \left[\frac{x}{h_{\Lambda}} \right] \right] = \left[\int_{X} \left[\frac{q^{s}_{\Lambda}}{h_{\Lambda}} \right] \right].$	Deterministic initial velocity field (0-discharge): $u(x, y, 0, \xi) = v(x, y, 0, \xi) = 0.$ Stochastic bottom topography:

using only the conserved variables.

- A sufficient condition to preserve the hyperbolicity, which is a stochastic variant of the deterministic positivity condition.
- A computationally tractable condition to guarantee the hyperbolicity.
- A central-upwind scheme that preserves both the hyperbolicity and the well-balanced property at discrete time levels.

Motivations

- Uncertainties can enter the shallow water system, for example, via the noisy measurement of the bottom.
- A SG formulation of shallow water equations is not necessarily hyperbolic.
- A non-well-balanced scheme may lead to spurious oscillations on relatively coarse grid.

Stochastic Parameterized Shallow Water System k=1• Stochastic Galerkin method applies standard Galerkin procedure in the stochastic ξ space, which leads to a new system of partial differential equations with respect to the PCE coefficients.



 $B(x, y, \xi) = 0.8e^{-5(x-0.9+0.1\xi^{(1)})^2 - 50(y-0.5+0.1\xi^{(2)})^2}.$ Randomness:

$$\xi^{(1)} \sim \text{Beta}(4,2), \qquad \xi^{(2)} \sim \mathcal{U}(-1,1).$$





The Main Results

Hyperbolicity-Preserving SG Formulation

$$\frac{\partial}{\partial t}(\widehat{U}) + \frac{\partial}{\partial x}(\widehat{F}(\widehat{U})) + \frac{\partial}{\partial y}(\widehat{G}(\widehat{U})) = \widehat{S}(\widehat{U},\widehat{B}).$$

Here,
$$\widehat{U} \coloneqq (\widehat{h}^{\top}, \widehat{q^{x}}^{\top}, \widehat{q^{y}}^{\top})^{\top}$$
, and

$$\widehat{F}(\widehat{U}) = \begin{pmatrix} \widehat{q^{x}} & & \\ \mathcal{P}(\widehat{q^{x}})\mathcal{P}^{-1}(\widehat{h})\widehat{q^{x}} + \frac{1}{2}g\mathcal{P}(\widehat{h})\widehat{h} \\ \mathcal{P}(\widehat{q^{y}})\mathcal{P}^{-1}(\widehat{h})\widehat{q^{y}} & \end{pmatrix}, \quad \widehat{G}(\widehat{U}) = \begin{pmatrix} \widehat{q^{y}} & & \\ \mathcal{P}(\widehat{q^{y}})\mathcal{P}^{-1}(\widehat{h})\widehat{q^{x}} & & \\ \mathcal{P}(\widehat{q^{y}})\mathcal{P}^{-1}(\widehat{h})\widehat{q^{y}} + \frac{1}{2}g\mathcal{P}(\widehat{h})\widehat{h} \end{pmatrix}, \quad \widehat{S}(\widehat{U}, \widehat{B}) = \begin{pmatrix} 0 \\ -g\mathcal{P}(\widehat{h})\widehat{B_{x}} \\ -g\mathcal{P}(\widehat{h})\widehat{B_{y}} \end{pmatrix}. \quad (\widehat{Q})$$

Theorem (Hyperbolicity-preserving condition)

The system (1) is hyperbolic if the matrix $\mathcal{P}(\hat{h}) > 0$.

The condition $\mathcal{P}(\hat{h}) > 0$ reduces to h > 0 when the ξ -dependence is dropped from the system.

$$\begin{split} &\frac{\partial h}{\partial t} + \frac{\partial q^x}{\partial x} + \frac{\partial q^y}{\partial y} = 0, \\ &\frac{\partial q^x}{\partial t} + \frac{\partial}{\partial x} \left(\frac{(q^x)^2}{h} + \frac{gh^2}{2} \right) + \frac{\partial}{\partial y} \left(\frac{q^x q^y}{h} \right) = -gh \frac{\partial B}{\partial x}, \\ &\frac{\partial q^y}{\partial t} + \frac{\partial}{\partial x} \left(\frac{q^x q^y}{h} \right) + \frac{\partial}{\partial y} \left(\frac{(q^y)^2}{h} + \frac{gh^2}{2} \right) = -gh \frac{\partial B}{\partial y}. \\ &h \text{ is the water height, } q_x \text{ and } q_y \text{ are the } x\text{- and } y\text{-} \\ &\text{ discharges, and } B \text{ is the time-independent surface.} \\ &\text{ All the variables are } \xi\text{-dependent random fields, e.g., } \\ &h = h(x, y, t, \xi). \end{split}$$

Polynomial Chaos Expansion (PCE)

• An unknown random field $z(x, y, t, \xi)$ is represented in the L^2_{ρ} orthonormal basis $\{\phi_k\}_{k \in \mathbb{R}}$, where $\rho \coloneqq \rho(\xi)$ is the density of the random parameter.

$$z(x, y, t, \xi) = \sum_{k=1}^{\infty} \widehat{z}_k(x, y, t)\phi_k(\xi),$$

• K-term truncated PCE:

Theorem (A computationally tractable condition)

Given Λ , let nodes ξ_m and weights τ_m satisfying $\{(\xi_m, \tau_m)\}_{m=1}^M$ represent any M-point positive quadrature rule that is exact on

 $P_{\Lambda}^3 \coloneqq \operatorname{span} \left\{ \prod_{n=1}^3 \phi_n \mid n \in [K] \right\}.$

 $h_{\Lambda}(x, y, t, \xi_m) > 0 \quad \forall \ m = 1, \dots, M,$

then the matrix $\mathcal{P}(\hat{h}) > 0$.

If

In other words, we only need to ensure the positivity of the stochastic water heights at some quadrature points to preserve the hyperbolicity of (1).

Second-Order Central-Upwind Scheme

Assuming uniform rectangular partition over a rectangular region, $\frac{d}{dt}\boldsymbol{U}_{i,j} = -\frac{\mathcal{F}_{i+\frac{1}{2},j} - \mathcal{F}_{i-\frac{1}{2},j}}{\Lambda r} - \frac{\mathcal{G}_{i,j+\frac{1}{2}} - \mathcal{G}_{i,j-\frac{1}{2}}}{\Lambda u} + \overline{\boldsymbol{S}}_{i,j},$ $\boldsymbol{U}_{i,j}^{E,W,N,S}$ are the pointwise values of the secondorder accurate, non-oscillatory piecewise linear reconstructions of $\boldsymbol{U}_{i,j}$ at the midpoints of the boundaries, i.e.,

(3)

$$\boldsymbol{U}_{i,j}^{E} = \boldsymbol{U}_{i,j} + \frac{\Delta x}{2} (\boldsymbol{U}_{x})_{i,j}, \quad \boldsymbol{U}_{i,j}^{W} = \boldsymbol{U}_{i,j} - \frac{\Delta x}{2} (\boldsymbol{U}_{x})_{i,j},$$
$$\boldsymbol{U}_{i,j}^{N} = \boldsymbol{U}_{i,j} + \frac{\Delta y}{2} (\boldsymbol{U}_{y})_{i,j}, \quad \boldsymbol{U}_{i,j}^{S} = \boldsymbol{U}_{i,j} - \frac{\Delta y}{2} (\boldsymbol{U}_{y})_{i,j},$$

Figure 1:Numerical results at T = 0.6 (top), T = 1.2 (middle), and T = 1.8 (bottom), respectively. The largest disks are corresponding to the standard deviation values 2.20e-3, 2.00e-3, and 1.20e-3, respectively. The index set $\Lambda = \{(\nu^{(1)}, \nu^{(2)}) \in$ $\mathbb{N}^2 \mid 0 \leq \nu^{(1)}, \nu^{(2)} \leq 3\}$. The polynomial basis is chosen to be the tensor-product set. Animation can be found in https:

 $\Pi_{\Lambda}[z]\coloneqq \sum_{k=1}^{K}\widehat{z}_{k}(x,y,t)\phi_{k}(\xi),$

where Λ is the index set for the (possibly multivariate) polynomials, the cardinality of Λ is K, and $\phi_1(\xi) = 1$.

• A K-term PCE approximation to product of two random fields a and b:

 $\Pi_{\Lambda}[a,b]\coloneqq\Pi_{\Lambda}[\Pi_{\Lambda}[a]\ \Pi_{\Lambda}[b]]$.

• A K-term PCE approximation to the ratio of two random fields a and b:

 $\Pi^{\dagger}_{\Lambda}[b/a]$: the solution to $\Pi_{\Lambda}[a,b/a] = \Pi_{\Lambda}[b].$

Notations

• $\hat{z} = (\hat{z}_1, \dots, \hat{z}_K)^\top$. • $\mathcal{P}(\hat{z}) \coloneqq \sum_{k=1}^K \hat{z}_k \mathcal{M}_k, \quad (\mathcal{M}_k)_{\ell m} = \langle \phi_k, \phi_\ell \phi_m \rangle_{\rho}$. • It can be shown that $\widehat{\Pi_{\Lambda}[a, b]} = \mathcal{P}(\hat{a})\hat{b}, \qquad \widehat{\Pi_{\Lambda}^{\dagger}[b/a]} = \mathcal{P}^{-1}(\hat{a})\hat{b}.$ where $U_{i,j}$ represent the cell averages of the vector \widehat{U} in rectangular cell $\mathcal{C}_{i,j}$.

• Source term:



 $\mathcal{G}_{i+\frac{1}{2},j} \coloneqq \frac{b_{i,j+\frac{1}{2}}^{+} \widehat{G}(\boldsymbol{U}_{i,j}^{N}) - b_{i,j+\frac{1}{2}}^{-} \widehat{G}(\boldsymbol{U}_{i,j+1}^{S})}{b_{i,j+\frac{1}{2}}^{+} - b_{i,j+\frac{1}{2}}^{-}} + \frac{b_{i,j+\frac{1}{2}}^{+} b_{i,j+\frac{1}{2}}^{-}}{b_{i,j+\frac{1}{2}}^{+} - b_{i,j+\frac{1}{2}}^{-}} \left[\boldsymbol{U}_{i,j+1}^{S} - \boldsymbol{U}_{i,j}^{N} \right].$



Hyperbolicity-Preserving Well-Balanced Central-Upwind Scheme

• Stochastic "lake-at-rest" state:

$$\begin{cases} q^x_{\Lambda} = q^y_{\Lambda} \equiv 0, \\ h_{\Lambda} + \Pi_{\Lambda}[B] \equiv C(\xi), \end{cases} \Rightarrow \begin{cases} \widehat{q^x} = \widehat{q^y} \equiv \mathbf{0}, \\ \widehat{h} + \widehat{B} \equiv \widehat{C}. \end{cases}$$

- The PCE vector B for the bottom function is replaced by its piecewise bilinear interpolant.
 The pointwise values of the reconstructions of the PCE of water surface η are reconstructed. The reconstructed water height are computed by h := η - B.
- The first moments \hat{h}_1 are "corrected" following a similar procedure to the central-upwind scheme for the deterministic shallow water equations.
- The PCE vectors \hat{h} are filtered to satisfies the condition (3).

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References

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